



Radio emission from ultra-diffuse galaxies residing in galaxy clusters

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Ultra-diffuse galaxies (UDGs), defined by their extremely low surface brightness (g -band $\mu \gtrsim 24$ mag arcsec⁻²) and large effective radii (3–10 arcsec), remain one of the most puzzling galaxy populations in the nearby Universe (van Dokkum et al., 2015). Predominantly found in dense environments, UDGs in the Coma cluster show a preferential alignment of their major axes toward the cluster centre, suggesting strong environmental influence on their formation and evolution. Using high-sensitivity, low-frequency radio data from the upgraded Giant Metrewave Radio Telescope (GMRT), a pathfinder instrument we examined all 854 UDGs cataloged in Coma cluster (Yagi et al., 2016). Despite the unprecedented depth of these observations, no individual detections were made. A median stacking analysis in the upgraded GMRT band-3 achieved a $5 \times \text{RMS}$ upper limit $\approx 1.5 \mu\text{Jy}$, providing the most stringent constraint yet on the average (mean) radio emission from UDGs, corresponding to star-formation rates $\lesssim 10^{-3} M_{\odot} \text{ yr}^{-1}$ for Coma-cluster-like systems and $\lesssim 10^{-1} M_{\odot} \text{ yr}^{-1}$ at $z \sim 0.05$.

Looking ahead, the Square Kilometre Array (SKA) will transform the study of such faint galaxies. While the early AA* configuration will deliver sensitivities comparable to the upgraded GMRT, the AA4 design baseline will achieve sub- μJy RMS levels at matched frequencies ($\nu \sim 200 \text{ MHz} - 1.4 \text{ GHz}$), enabling detections of UDGs with star formation rates as low as $10^{-4} - 10^{-3} M_{\odot} \text{ yr}^{-1}$ within Virgo and Coma distances. Such capabilities will allow robust discrimination between quenched, dark-matter-dominated systems and those sustaining weak residual star formation or low-luminosity nuclear activity.

Clearly, the forthcoming sensitivity of SKA-Mid and SKA-Low will be decisive in determining whether UDGs extend the faint tail of the active galaxy or star-formation luminosity function, or instead constitute a distinct radio-silent class. Even if the future looks dark for UDGs, in the sense that their radio emission may remain extremely challenging to detect, there is still good hope that upcoming facilities will possibly reveal their true nature.

1 Introduction

Ultra-diffuse galaxies (UDGs) are an extreme population of low-surface-brightness systems, especially abundant in rich galaxy clusters. [van Dokkum et al. \(2015\)](#) defined them as galaxies with central g -band surface brightness $\mu_g \gtrsim 24$ mag arcsec⁻² and large effective (half-light) radii $R_e \gtrsim 1.5$ kpc (see [Van Nest et al., 2022](#), for an in-depth discussion on UDG definition). Thousands of UDGs have now been identified across a wide range of environments, e.g., Abell 2744: [Janssens et al. \(2019\)](#), Coma cluster: [Yagi et al. \(2016\)](#), [Zaritsky et al. \(2019, 2021\)](#), Globular Cluster system: [Gannon et al. \(2024\)](#), Pisces-Perseus super-cluster: [Martínez-Delgado et al. \(2016\)](#), Perseus cluster: [Gannon et al. \(2022\)](#), Stripe 82: [Barbosa et al. \(2020\)](#), UGC 6594 group: [Gannon et al. \(2021\)](#), Virgo cluster: [Mihos et al. \(2015\)](#), etc. Observed cluster UDGs span a substantial luminosity range, an absolute magnitude of about -12 to -16 ([Koda et al., 2015](#)), although the completeness at the faint end of their luminosity function remains uncertain. Recent studies suggest that more than 7% of all galaxies may be ultra-diffuse ([Li et al., 2023](#)). Understanding how these galaxies form and evolve has therefore become a challenge in galaxy formation research, with UDGs emerging as valuable laboratories for studying both galaxy formation and cluster evolution.

A wide variety of theories have been proposed to explain the formation of UDGs, generally invoking either (i) external processes such as tidal heating, ram-pressure stripping, environmental quenching, or galaxy mergers ([Carleton et al., 2019](#); [Sales et al., 2020](#); [Jones et al., 2021](#); [van Dokkum et al., 2022](#), etc.), or (ii) internal mechanisms such as high dark-matter halo spin, stellar feedback, or passive stellar evolution ([Di Cintio et al., 2017](#); [Benavides et al., 2023](#); [Fielder et al., 2024](#), etc.). In many cases, a combination of both internal and external factors possibly contributes to their formation (see also [Martin et al., 2019](#); [Gannon et al., 2024](#), and references therein), and different formation mechanisms are expected to leave distinct signatures in the stellar populations and dark matter halos of UDGs. For instance, UDGs formed through episodic stellar feedback should exhibit extended star formation histories, low metallicities, and dwarf-galaxy-like halos with correspondingly low velocity dispersions and few globular clusters ([Zheng et al., 2024](#)). In contrast, UDGs that formed early and quenched rapidly are expected to host old, metal-poor stellar populations and reside in more massive halos, with higher velocity dispersions and richer globular-cluster-like systems ([Gannon et al., 2024](#)).

Observational studies have shown a transition in the overall properties of field and cluster UDGs, with UDGs in the field tending to be bluer, more irregular, and gas-rich, with some ongoing star formation, while those in clusters are mostly red, spheroidal, and quenched (e.g. [van der Burg et al., 2016](#); [Román and Trujillo, 2017a,b](#)). Although most field UDGs are gas-rich and star-forming, a population of isolated quenched systems has also been reported, which may originate as backslash galaxies that were previously satellites of groups or clusters, as suggested by simulations ([Benavides et al., 2021](#)). Large optical surveys and dedicated catalogs (e.g., [Yagi et al., 2016](#)) have revealed that clusters such as Coma host hundreds of UDGs, far exceeding field expectations, and that their structural properties (e.g. axis ratios, sizes) are remarkably coherent, suggesting that environmental effects play a significant role in shaping them ([Motiwala et al., 2025](#)). Notably, (i) many UDGs appear as barely resolved diffuse systems along the line of sight to the Coma cluster ($R_e \sim 800$ pc—5 kpc, effective surface brightnesses = 25–28 mag arcsec⁻² and stellar masses $\sim 1 \times 10^7 M_\odot$ –

$5 \times 10^8 M_{\odot}$, Koda et al., 2015), and (ii) deep stacked Subaru images show little evidence of strong tidal disturbances in their outskirts, implying either substantial dark-matter dominance or survival mechanisms that shield them from rapid disruption (Yagi et al., 2016).

Results from cosmological simulations, e.g., Millennium-II and Phoenix: Rong et al. (2017), RomulusC: Tremmel et al. (2020), TNG50: Benavides et al. (2023), EAGLE: Zheng et al. (2025), NIHAO: Motiwala et al. (2025), etc., present a picture in which UDGs are not a wholly separate galaxy class, but rather the extreme tail of the dwarf-galaxy population emerging under particular conditions. Many simulated UDGs occupy dwarf-galaxy-like mass halos ($M_{\star} \sim 10^7\text{--}10^9 M_{\odot}$), and yet are unusually extended (effective radii $\gtrsim 2\text{--}3$ kpc) because of high halo spin, early quenching, or tidal/heating processes. In simulations focusing on clusters of galaxies (van Dokkum et al., 2022), UDG analogues often fall into massive halos early and suffer strong environmental processing, e.g., ram-pressure stripping, tidal mass loss and adiabatic expansion, leading to low surface brightness and gas-poor states (Tremmel et al., 2020). Alternatively, in field environments, simulations show a subset of “born UDGs” whose gas spins and feedback-driven expansions yield large radii without heavy stripping. Across environments the result is that UDGs tend to have large sizes for their stellar mass, older stellar populations, and a continuous rather than discrete distribution with respect to normal dwarf galaxies (Benavides et al., 2023). Thus, understanding UDGs in clusters is important because their origin, whether they are “puffed-up” dwarfs, failed milky-way-like galaxies, or a mix of both, which contains key information about galaxy formation at low surface brightness and the environmental processes that quench star formation. Their abundance scales with the mass of their host halo, makes them valuable tracers of dark-matter-dominated regions and environmental processes during the formation and growth of galaxy clusters that affect their member galaxies.

Although optical surveys have established large samples of UDGs across a range of environments, only a few attempts have been made to detect them at radio wavelengths. Early efforts targeted their HI and continuum emission to probe gas content and star formation activity. Papastergis et al. (2017) used the Effelsberg 100 m telescope to search for HI in four isolated UDGs, detecting one (SdI-2) with $M_{\text{HI}} \approx 2.4 \times 10^8 M_{\odot}$, while setting upper limits of $(1.3\text{--}2.4) \times 10^8 M_{\odot}$ for the others. Scott et al. (2021) presented resolved HI observations of two blue, gas-rich (isolated) UDGs using the GMRT, confirming that such systems exist outside clusters but differ from the quiescent, red UDGs found in dense environments. In contrast, Struble (2018) searched for radio counterparts to UDGs in the Coma cluster using VLA data and found no significant matches above background levels, indicating that cluster UDGs are largely radio quiet. More recently, the SMUDGes HI survey (Karunakaran et al., 2024) reported detections for 110 out of 378 UDG candidates, predominantly in isolated environments, again confirming that gas-rich UDGs are primarily field systems while cluster UDGs are strongly quenched. This motivates deep, targeted radio continuum and HI observations of UDGs using current and upcoming facilities. Pathfinder and precursor arrays (upgraded GMRT, MeerKAT, LOFAR), with their exceptional sensitivity and wide-field imaging capabilities, are ideally suited to bridge the gap between deep optical surveys and the unprecedented resolution and sensitivity that the SKA will deliver. Such observations will enable the first statistically significant constraints on the non-thermal emission, magnetic fields, and cold-gas reservoirs of UDGs across a range of environments.

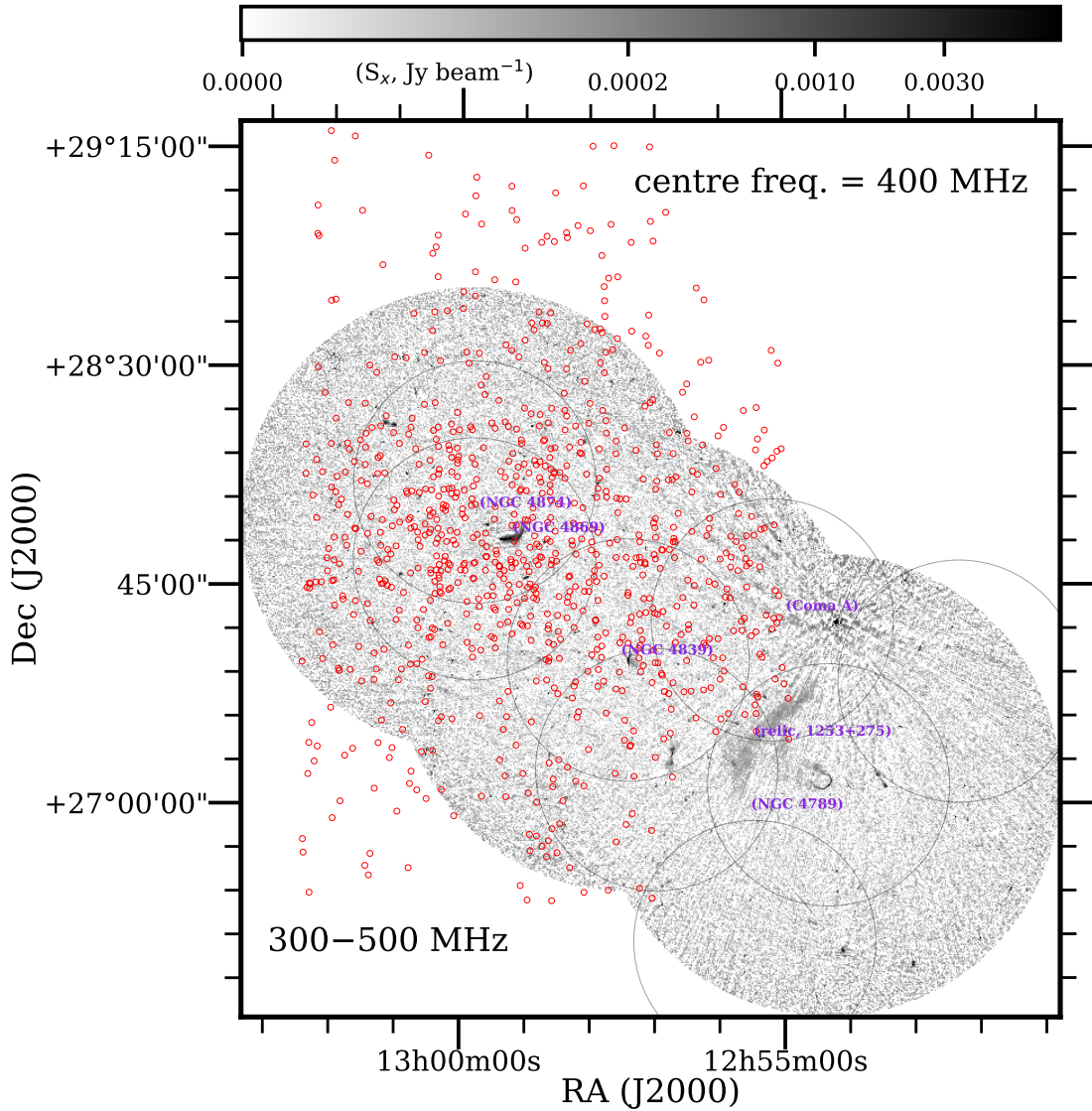


Figure 1: Mosaic image of three pointings in the upgraded GMRT band-3 (250–500 MHz) at an angular resolution of $\sim 6.5''$. UDG positions from Yagi et al. (2016) are marked by the “o” sign. The gray-scale image is displayed in logarithmic scales to emphasize the extended, low-surface brightness diffuse radio emission, and the names of key sources are labeled. The rms noise is relatively uniform between $21 \mu\text{Jy beam}^{-1}$ and $31 \mu\text{Jy beam}^{-1}$, apart from small regions near bright sources, e.g., Coma A where the rms noise increases. Eight circles correspond to the eight pointings in the 550–850 MHz band of upgraded GMRT. Note the starry-pattern artifact centered on Coma A, limiting the dynamic range.

2 First results from upgraded GMRT

Although pathfinder and precursor facilities have achieved remarkable sensitivity to diffuse cluster emission, direct radio detections of individual UDGs remain rare. The UDG sample is drawn from the catalog of Yagi et al. (2016), who identified 854 such systems in the Coma cluster ($z = 0.02316$) through a systematic search of deep R -band Suprime-Cam/Subaru imaging. Their study provides a

comprehensive and homogeneous compilation of UDGs in a dense cluster environment, forming a well-defined sample for statistical investigations. We use this optically selected catalog as the basis for a radio study, aiming to examine the presence and properties of faint radio emission from these galaxies and to explore their contribution to the low-luminosity end of the radio population.

We therefore analyzed the Coma cluster UDGs using archival low-frequency radio continuum observations obtained with the upgraded Giant Metrewave Radio Telescope (GMRT), a pathfinder instrument. The UDGs lie within the area covered by the deep upgraded GMRT survey of the Coma cluster (Lal et al., 2022). We use observations in both band 3 (250–500 MHz) and band 4 (550–850 MHz), which together provide high-sensitivity, arcsecond-scale resolution across the cluster environment. These final images have a RMS noise level of 21.1–36.6 $\mu\text{Jy beam}^{-1}$ and 12.8–42.4 $\mu\text{Jy beam}^{-1}$ at the angular resolutions of $\sim 6.1''$ and $\sim 3.7''$ in the 250–500 MHz and 550–850 MHz bands, respectively. After matching the two images at an angular resolution of $\sim 6.5''$, we extracted a region of $50'' \times 50''$ from these band-3 and band-4 images centered on the UDG position as given by the Subaru survey for each UDG in our sample. We retained 736 and 582 UDGs suitable for stacking in band 3 and band 4, respectively. Figure 1 shows the mosaic image of three pointings in the upgraded GMRT band-3 (250–500 MHz) at an angular resolution of $\sim 6.5''$. Eight circles correspond to the eight pointings in the 550–850 MHz band of upgraded GMRT. UDG positions are marked by the “ \circ ” sign.

Surprisingly, no individual UDG was detected in either band above the 3σ ($= 3 \times \text{RMS}$) threshold within the sensitivity limits of the images. Thus, we explored two complementary methods for stacking radio images, mean and median stacking, each with distinct advantages and limitations (White et al., 2007). Mean stacking provides an easily interpretable measure of average flux density but is highly sensitive to outliers, such as bright sources or noisy images within the stack, which can bias the result. Although such effects can be mitigated by excluding sources above certain flux density or RMS thresholds within the extracted region, the result remains sensitive to the specific cutoff criteria. Median stacking, on the other hand, is considerably more resilient to outliers and non-Gaussian noise distributions (e.g., Gott et al. 2001), allowing all images to be included without arbitrary truncation. However, at low signal-to-noise ratios, interpreting the median becomes less straightforward because the recovered value can be shifted toward the local mean, with the magnitude of the shift depending on the noise level. Thus, while the median provides stability against contamination, it can underestimate the true flux density for faint noise-dominated samples.

To investigate the ensemble radio properties of these ultra diffuse galaxies, we performed mean and median stacking analyses in both bands, which offers a valuable cross-check, ensuring a more reliable characterization of the statistical radio properties. In band 3, we stacked 675 regions with no outliers for mean-stacking and all 736 images for median-stacking, whereas in band-4, we stacked 533 regions with no outliers for mean-stacking and all 582 images for median-stacking. We did not detect any signal at the locations of the UDGs in any of the stacked images. In band-3, the stacked image yielded an average flux density upper limit of $\sim 2.3 \mu\text{Jy}$ for mean-stacking and $\sim 1.7 \mu\text{Jy}$ median-stacking. Whereas in band 4, the stacked image yielded an average flux density upper limit of $\sim 1.6 \mu\text{Jy}$ for mean-stacking and $\sim 1.5 \mu\text{Jy}$ median-stacking. Figure 2 shows resulting image at band-3 ($= 50'' \times 50''$) of constructing a mean (left panel) and a median (right panel) stack of

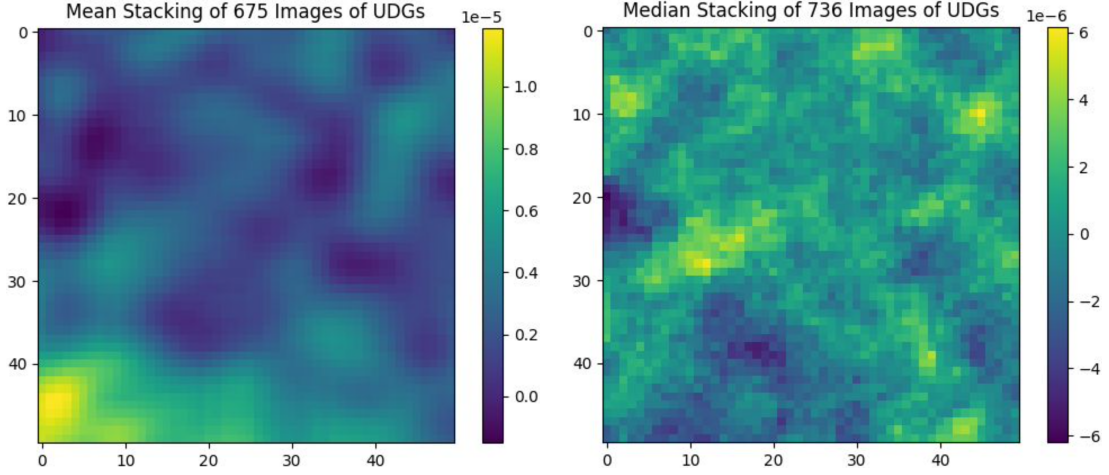


Figure 2: Result of constructing a mean (left panel) and a median (right panel) stack of the 675 and 736, respectively source positions in the UDG catalog. The image displays a $50'' \times 50''$ color-scale image (pixel size = $0.8''$). The flux densities from the stacked images are consistent with the expected noise reduction ($\approx N^{-1/2}$), where N is the number of images in the stack, confirming that the results are dominated by Gaussian noise rather than faint unresolved sources.

the 675 and 736, respectively source positions in the UDG catalog. The resulting band-4 stacked images are (again non-detection and are therefore) not shown. The flux densities measured from the stacked images are consistent with the expected noise reduction ($\approx N^{-1/2}$, where N is the number of images in the stack), confirming that the results are dominated by Gaussian noise rather than faint unresolved sources. We have also analyzed LOFAR LoTSS DR2 data (Shimwell et al., 2022) of the Coma cluster, but the prominent diffuse emission from the Coma radio halo limits our ability to reach deeper sensitivities through stacking.

We note that scaling the band-3 (central frequency = 400 MHz) limits assuming a spectral index of $\alpha = -0.7$ yields a limit of $\sim 1.1 \mu\text{Jy}$ at band-4 (central frequency = 700 MHz), which is about a factor of 1.4 higher than the current best estimates ($\approx 1.5 \mu\text{Jy}$) derived from the stacked images. Thus, our median-stacked image at band-3 is the deepest among all data-sets, providing the most stringent observational limits on the average radio emission from UDGs. Consequently, we adopt this stack for all subsequent analyses. The combination of broader bandwidth, better sensitivity, and a larger number of contributing sources makes the band-3 result particularly robust for constraining faint, diffuse synchrotron emission in these ultra-diffuse systems.

Upper limit on the star-formation rate Most of the non-thermal radio emission from galaxies arises from electrons that have diffused away from their acceleration sites in supernova remnants into quiescent regions void of current star formation, and using Murphy et al. (Eq. 17, 2011), we determine star formation rate (SFR),

$$\text{SFR (M}_\odot \text{ yr}^{-1}) = \frac{L_{400} \left(\frac{1.4 \text{ GHz}}{0.4 \text{ GHz}} \right)^\alpha}{1.57 \times 10^{21} \text{ W Hz}^{-1}}, \text{ and}$$

$$L_\nu \text{ (W Hz}^{-1}) = 4\pi D_L^2 (S_\nu \text{ (}\mu\text{Jy)} \times 10^{-32});$$

where D_L is the luminosity distance ($= 3.086 \times 10^{24}$ m), and we assume $\alpha = -0.7$. These imply that the mean SFR of UDGs in Coma cluster is $\approx 4.8 \times 10^{-4} M_\odot \text{ yr}^{-1}$, well below $10^{-3} M_\odot \text{ yr}^{-1}$, consistent with their quiescent optical colors and lack of individual radio counterparts. Together, these findings support that UDGs in rich clusters are strongly quenched systems, possibly shaped by the dense intra-cluster medium and environmental stripping processes. At present, these constitute the best observational limits on our understanding of UDGs. This motivates deep, targeted radio continuum and HI observations, especially with the SKA, to constrain their cold-gas reservoirs, star-formation and active galactic nuclei activity, and to probe how cluster environments shape and regulate the formation and survival mechanisms of these extremely diffuse galaxies.

3 Predictions for SKA AA* and AA4

We note that the AA4 (design baseline) configuration forms the reference (Bourke et al., 2015; Braun, 2015). Specifically, the SKA-Mid array comprises 197 dishes, 133 15-m SKA dishes and 64 13.5-m MeerKAT dishes (with baselines out to 150 km), while the SKA-Low array consists of 512 low-frequency aperture array stations of 256 antennas each (with a maximum baseline of around 75 km). On the other hand, AA* is a reduced-scale, interim deployment by the SKA Observatory (Braun et al., 2024, and talk by Shari Breen, Head of Science Operations, 2025 October 20, “Meeting with the SKAO and Book Chapters”). Despite its smaller collecting area and correspondingly lower sensitivity compared to AA4, AA* will support the full range of observing modes, e.g., up to 16 subarrays, commensal operations, and a wide frequency coverage from 50 MHz to 15.4 GHz. In this configuration, SKA-Mid includes 144 dishes (with a maximum baseline of 36 km), and SKA-Low consists of 307 stations (with a maximum baseline of 74 km).

Expected RMS sensitivities Thus, the interim AA* configuration and later the baseline AA4 configuration of the SKA Observatory represent a significant step forward in addressing these questions. With continuum sensitivities at 1.4 GHz and 200 MHz exceeding those of the upgraded GMRT’s band-3 (and band-4), AA4 and AA* will be capable of probing average flux densities of order $0.18 \mu\text{Jy}$ and $0.21 \mu\text{Jy}$, and $0.60 \mu\text{Jy}$ and $0.78 \mu\text{Jy}$ in long ($= 100$ hr) integrations for SKA-Mid and SKA-Low respectively, which is roughly an order of magnitude deeper than current stacking limits ($\sim 1.5 \mu\text{Jy}$). At these levels, it will be possible to distinguish between thermal star-formation-driven emission and weak non-thermal synchrotron signatures from embedded active galaxy, even in the faintest cluster UDGs. Moreover, the combination of wide field-of-view, excellent (u, v)-coverage, and arcsecond-scale resolution will enable both statistical studies of large UDG samples and the separation of UDG emission from surrounding diffuse cluster radio halo emission.

3.1 Expected star-formation rate

We assume (i) point-source emission, (ii) a spectral index, $\alpha = -0.7$, (iii) the adopted radio-SFR calibration, and (iv) no significant loss of flux density due to resolved-out low-surface-brightness emission. Clearly, realistic estimates should thus include spectral-index uncertainties, account for effects to the SFR due to cosmic-ray, i.e., systems can under-produce synchrotron emission for a given SFR due to cosmic-ray, and ability to image extended sources faithfully due to limited (u, v)-coverage of radio interferometers.

Figure 3: The estimates of 5σ detection limits and star-formation rate sensitivities for SKA AA4 and AA* configurations.

Telescope	Config.	Frequency (GHz)	Integration time (hr)	RMS 5σ Flux density (μJy)	Equiv. flux density 1.4 GHz	SFR $_{5\sigma}$			
						Virgo ($z = 0.00436$)	Coma ($z = 0.02316$)	$z = 0.05$	
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)
SKA-Mid	AA4 ^a	1.4	8 [†]	0.24	1.20	1.20	0.25	9.14	44.20
			100	0.07	0.35	0.35	0.07	2.67	12.90
	AA* ^b		8 [†]	0.34	1.70	1.70	0.35	13.00	62.60
			100	0.10	0.50	0.50	0.10	3.81	18.40
SKA-Low	AA4 ^c	0.2	8	39.44	197.20	50.47	10.50	385.00	1860.00
			100	11.85	59.25	15.17	3.16	116.00	558.00
	AA* ^d		8	63.96	319.80	81.87	17.00	624.00	3010.00
			100	19.21	96.05	24.58	5.12	187.00	905.00

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We use $A_{\text{eff}}/T_{\text{sys}}$ for SKA1-Mid from SKAO-TEL-0000818 (Table in Appendix B, Braun et al., 2024) at central frequency = 1.396 GHz (≈ 1.4 GHz), representative continuum bandwidth = 770 MHz (typical SKA-Mid instantaneous usable bandwidth around L-band), no. of polarisations = 2, and natural weighting (see also Fanaroff et al., 2021, as published use case). Below we give system-equivalent flux density (SEFD), and continuum sensitivity (in col. 5) determined using the radiometric equation for the SKA-Mid.

^a $A_{\text{eff}}/T_{\text{sys}} = 1751.0 \text{ m}^2 \text{ K}^{-1}$, gives SEFD = 1.577 Jy.

^b $A_{\text{eff}}/T_{\text{sys}} = 1216.2 \text{ m}^2 \text{ K}^{-1}$, gives SEFD = 2.270 Jy.

We use central frequency = 200 MHz, continuum bandwidth = 300 MHz (no. of sub-bands = 1), weighting correction factor = 30% continuum bandwidth and uniform weighting. (We do not use natural weighting here, since it results in poor fits due to large dirty-beam sidelobe levels.) Below we give weighted continuum sensitivity (in col. 5) determined using exposure times calculated with the SKA Sensitivity Calculator at <https://sensitivity-calculator.skao.int> for the SKA-Low.

^c Using 512 stations, synthesized beam-size = $3.89'' \times 3.03''$. The confusion noise, $\sigma_c = 1.04 \mu\text{Jy beam}^{-1}$.

^d Using 307 stations, synthesized beam-size = $3.82'' \times 3.44''$. The confusion noise, $\sigma_c = 1.21 \mu\text{Jy beam}^{-1}$.

[†] We note that in an 8 hr long integration time using 58 dishes (with a maximum baseline of 8 km) of the MeerKAT, the expected thermal noise sensitivity for continuum observations in the L-band (881–1670 MHz) = $3.2 \mu\text{Jy beam}^{-1}$, assuming SEFD = 425 Jy, and effective available bandwidth = 385 MHz; from Table 1, available at <https://skaafrica.atlasiian.net/wiki/spaces/ESDKB/pages/1486750321/Sensitivity+calculators>. Scaling this to 144 dishes (with a maximum baseline of 36 km) and 197 dishes (with a maximum baseline of 150 km), gives thermal noise sensitivities of $1.28 \mu\text{Jy beam}^{-1}$ and $0.94 \mu\text{Jy beam}^{-1}$, respectively for AA* (interim) and AA4 (design baseline). Although these estimates are a factor of ~ 4 higher than our estimates above (col. 5), we believe these are due to improved $A_{\text{eff}}/T_{\text{sys}}$ and hence improved specifications, $A_{\text{eff}}/T_{\text{sys}}$ (or SEFD) of the SKA-Mid array. Additionally, the confusion noise ($\sigma_c \propto \theta^2 \alpha^{-0.7}$, Condon et al., 1998) would be lower by 2–3 orders of magnitude (or lower by a factor of 40–700).

Table 3 summarises the 5σ detection limits and corresponding SFR sensitivities for the SKA-Mid and SKA-Low AA4 and AA* configurations for integration times of 8 hr and 100 hr. These predictions indicate that the SKA will be capable of detecting extremely low levels of star formation, on the order of $\sim 10^{-6} M_{\odot} \text{ yr}^{-1}$ in UDGs within the Virgo cluster ($z = 0.00436$), and $\sim 10^{-3} - 10^{-1} M_{\odot} \text{ yr}^{-1}$ for UDGs at $z \approx 0.05$. With these sensitivities, both AA* and especially AA4 should be able to detect the integrated synchrotron emission from UDGs at $z \approx 0.05$ and beyond. Instead, if such galaxies are truly quenched, the resulting non-detections will yield stringent upper limits on their radio emission, providing key constraints on feedback processes and quenching mechanisms in dense environments. Even if the future looks dark for UDGs, there is still good hope that upcoming facilities, SKA will reveal their true nature. Hence, AA4 observations are a must to test whether UDGs are predominantly dark-matter-dominated, quiescent systems or whether residual star formation and weak active galaxy still persist in cluster environments.

4 Discussion

Deep optical surveys of nearby clusters (e.g., Coma, Virgo: Yagi et al., 2016; Mihos et al., 2015) find that UDGs occupy the faint, largest, most extended members of the dwarf galaxy population, and that clusters host hundreds of UDGs, many more UDGs than would be predicted from the field population, even after correcting for observational incompleteness and selection biases. When UDGs in cluster environments are included, they can steepen the measured faint-end of the cluster galaxy luminosity function; note that the exact slope and normalization at absolute magnitude $\gtrsim -12$ remain uncertain because most (spectroscopic) surveys are limited by a lack of understanding as to the selection biases, e.g., surface-brightness completeness, effective (half-light) size cut, area coverage, etc. Cosmological and zoom simulations predict that the faint end of the UDG abundance is strongly sensitive to host-halo mass, infall time, and feedback history, and that the lowest-mass systems are strongly affected by tidal effects. In short, the simulations imply that UDGs in clusters are largely quenched, extended dwarfs that are formed via dwarf-environment interactions, while field UDGs may remain gas-rich but diffuse and not centrally concentrated, i.e., low-surface-density. Such theoretical frameworks help interpret the extremely faint radio and star-formation signatures of UDGs as natural outcomes of their formation scenarios.

UDG abundance Thus, UDGs have emerged as a key population for probing the link between galaxy formation and environment, particularly within groups and clusters where their number found in a system tends to increase as the mass of the system’s dark-matter halo increases. Early studies suggested that richer environments contain disproportionately more UDGs, hinting at a possible environmental origin or enhanced survival in massive haloes. Recent work by Makda et al. (2025), using a sample of 51 Hyper Suprime-Cam Subaru identified clusters spanning halo masses from $M_{200} \simeq 0.9 \times 10^{14} - 8.3 \times 10^{14} M_{\odot}$, finds that the abundance of UDGs follows a relatively shallow scaling, $N_{\text{UDG}} \propto M_{200}^{0.78 \pm 0.28}$. This slope is noticeably flatter than the steeper dependence ($N_{\text{UDG}} \propto M_{200}^{1.11 \pm 0.07}$) reported by van der Burg et al. (2017) based on GAMA groups and clusters. Over the same mass range, the Makda et al. (2025) relation predicts an increase in UDG counts by a factor of about 5 from low- to high-mass clusters, whereas the steeper van der Burg et al. (2017) relation predicts nearly twice that rise (~ 11). In other words, the shallower slope

may suggest that UDG formation or survival efficiency saturates toward massive haloes, possibly reflecting enhanced disruption, tidal destruction, or completeness limitations at the faint end, and the steeper slope suggest that dense cluster environments promote UDG survival or production more effectively. These contrasting slopes therefore have significant implications for interpreting the UDG population, determining whether they trace halo mass in proportion or instead represent a population whose abundance plateaus in the most massive clusters.

Presently, it seems that UDGs represent an extension of the extremely faint, low-luminosity active galactic nucleus population (Ho et al., 1997; Ulvestad and Ho, 2001, 2002; Lal et al., 2011) observed in nearby galaxies. Therefore, similar to low-ionisation nuclear emission-line regions (LINERs) and other weakly active nuclei Ulvestad and Ho (2001, 2002), UDGs could host accretion processes operating at very low Eddington ratios ($\lesssim 10^{-5}$ – $10^{-3} L_{\text{Edd}}$), where radiatively inefficient accretion flows and weak or absent jets produce minimal radio output (Lal and Ho, 2010; Lal et al., 2011). The absence of detectable radio emission from stacked UDG samples is consistent with the faint end of the active galaxy luminosity function (Bontempi et al., 2012), since radio emission may be suppressed by free–free absorption, where the absorber is typically ionised gas ($n_e = 10^{-5}$ – 10^{-3} cm^{-3} , $T = 10^4 \text{ K}$) in the inner narrow-line region or circumnuclear disk, or simply fall below the sensitivity of current instruments. If some UDGs indeed host faint, compact, low-power active galaxy (e.g., Baldi et al., 2019), their detection would provide valuable constraints on black-hole occupation fractions and feedback efficiency in low-mass, diffuse systems.

Note that dust absorption is negligible in the radio band but can play a significant role in the infrared, particularly for UDGs with low-level star formation (see also Magliocchetti et al., 2025). Translating the inferred star-formation rate limits from our radio stacking ($\sim 10^{-3}$ – $10^{-2} M_{\odot} \text{ yr}^{-1}$) into expected infrared fluxes using the radio–infrared correlation suggests that such UDGs would emit at levels of a few μJy at 100–250 μm . At the distances of Virgo (16.5 Mpc) and Coma (100 Mpc), this corresponds to dust temperatures of roughly 15–25 K, assuming standard blackbody models, and places them well below current Euclid (Euclid Collaboration et al., 2025) and Herschel (Pilbratt et al., 2010) detection limits. For more distant systems ($z \sim 0.05$), the corresponding fluxes would still be fainter ($< 1 \mu\text{Jy}$), implying that any dust present in these galaxies must be extremely cold and of very low mass. Thus, even next-generation infrared surveys will struggle to detect the dust emission from such systems unless their star formation is episodically enhanced.

4.1 Looking ahead

Clearly, UDGs occupy an extreme regime: large effective radii combined with low stellar densities challenge standard models of galaxy formation and feedback. Some UDGs host unexpectedly massive dark matter halos or rich globular cluster systems, indicating diverse formation scenarios. Their low surface brightness and tenuous interstellar medium make them highly sensitive probes of gas stripping, quenching, and star-formation efficiency in dense environments. Table 1 summarizes the likely physical causes for the extremely weak or undetected radio emission observed in UDGs and in galaxies at the faint end of the radio luminosity function, such as LINERs, H II region galaxies, etc. The absence of significant synchrotron radiation in these systems may reflect intrinsically weak magnetic fields and inefficient cosmic-ray acceleration, rapid electron energy losses, and free–free absorption or synchrotron self-absorption. Environmental effects such as ram-pressure stripping

Table 1: Summary of possible causes for extremely weak or undetected radio emission from galaxies at the faint end of the luminosity function, including UDGs and low-luminosity active galaxies.

Possible cause	Physical explanation and implications
1. INTRINSICALLY WEAK SYNCHROTRON EMISSION	Low cosmic-ray electron densities or inefficient particle acceleration in weak active galaxy jets and low star-formation environments lead to faint non-thermal emission. In UDGs, both magnetic field strengths ($B \lesssim 1\text{--}2 \mu\text{G}$) and cosmic-ray densities are likely sub-normal (Hardcastle et al., 2026).
2. DOMINANCE OF THERMAL (FREE-FREE) EMISSION WITH LOW NORMALIZATION	In galaxies with extremely low star formation rates ($\lesssim 10^{-3}\text{--}10^{-2} M_{\odot} \text{yr}^{-1}$), thermal radio emission can dominate but remains below current sensitivity thresholds, yielding no measurable spectral signature (An et al., 2026).
3. ENERGY LOSSES OF COSMIC-RAY ELECTRONS	Rapid inverse Compton losses against the cosmic microwave background (especially at higher redshift) and synchrotron cooling in weak magnetic fields can suppress radio flux density at GHz frequencies. The effect is particularly important for diffuse, low-B systems such as UDGs and LINER-like nuclei.
4. FREE-FREE ABSORPTION IN COMPACT STAR-FORMING OR ACTIVE GALAXIES	At low frequencies ($\lesssim 300$ MHz), partially ionised gas (e.g., narrow-line region, $T \sim 10^4$ K) can absorb synchrotron emission. Although dust absorption is negligible in radio bands, compact ionised gas or circumnuclear material could obscure low-level radio cores. Explains many LINER/H II and radio-quiet sources (Algera et al., 2026; An et al., 2026).
5. LOW ACCRETION-RATE (RADIATIVELY INEFFICIENT) ACTIVE GALAXIES	Many faint LINERs and transition-type nuclei, e.g., H II galaxies, radio-quiet sources host radiatively inefficient accretion flows (or advection dominated accretion flows), which produce weak or no observable radio jets. Such mechanisms could explain the absence of radio emission in low-luminosity active galaxies analogs among UDGs.
6. ENVIRONMENTAL GAS STRIPPING AND FEEDBACK	Cluster UDGs likely experienced strong ram-pressure stripping and tidal interactions that removed cold gas reservoirs, halting both star formation and active galactic nuclei fueling, thereby quenching radio emission.
7. OBSERVATIONAL SENSITIVITY LIMITS AND SURFACE-BRIGHTNESS DILUTION	For extended, low-surface-brightness galaxies, flux density may be spread over large angular scales, falling below detection thresholds. Deep stacking (median or mean) can recover integrated limits, but single detections remain challenging even with current high-sensitivity instruments (Mazzolari et al., 2026).
8. INTERMITTENT / EPISODIC OR STOCHASTIC ACTIVITY CYCLES	Short-lived, recurrent radio episodes in low-luminosity active galaxies may make detection probability small in snapshot surveys. SKA time-domain studies could test this hypothesis.
Towards next-generation (SKA) observations	The SKA (AA4 configuration) will deliver μJy - to nJy -level sensitivities, enabling direct detections or stringent upper limits for UDGs and faint LINERs out to $z \sim 0.05$. Combining SKA continuum and HI data will separate gas-rich, star-forming UDGs from quenched, dark systems, constraining the interplay between gas, feedback, and weak radio emission from the galaxy nucleus.

further quench star formation and active galactic nuclei fueling. Dust plays little role at radio wavelengths, but ionised gas absorption, the low surface brightness extended diffuse emission, and transient low-power active galaxy activity may also limit detection at radio wavelengths. In some cases, radiatively inefficient accretion or free–free absorption within compact ionised regions may further suppress any possible radio emission. Hence, multi-wavelength studies, especially deep radio continuum and H I observations are crucial for revealing residual star formation, cosmic-ray activity, and gas content, providing complementary constraints to optical surveys.

Clearly, UDGs bridge the gap between weak dwarf systems and powerful galaxies, trace the faint end of the galaxy luminosity function, and provide stringent tests of both astrophysical and cosmological models. By studying UDGs, we gain not only a more complete census of galaxies in clusters but also a deeper understanding of the physical processes shaping the low-mass, low-density universe. Their ubiquity, extreme properties, and sensitivity to environment make them indispensable to our understanding of cosmic structure and evolution.

Future studies with the SKA, promise to uncover the faintest regimes of UDG activity The cosmological simulations predict clear systematic trends in properties of UDGs. For example, (i) UDGs typically reside in dark-matter halos of $10^{10-11} M_{\odot}$, corresponding to high mass-to-light ratios (up to $\sim 10^3$ within R_e), and show quenched or weakly star-forming stellar populations by $z = 0$. (ii) In high-resolution models, the star-formation rates rarely exceed 10^{-3} – $10^{-2} M_{\odot} \text{ yr}^{-1}$, consistent with a near-quiescent state, and feedback-limited radio flux densities (or infrared fluxes) below current detection thresholds. (iii) The predicted gas fractions decline sharply with increasing environmental density, while rotation-dominated kinematics are more common in field UDGs than in cluster UDGs, which often exhibit inflated velocity dispersions and disturbed morphologies. Moreover, simulated luminosity functions show a flattening of the faint-end slope ($\alpha \approx -1.1$ to -1.3) in dense clusters compared to the field ($\alpha \approx -1.5$), reflecting strong environmental quenching. However, despite recent progress in observations using best, SKA pathfinder and precursor facilities, several key questions, listed below, remain unresolved for UDGs and can be addressed with the SKA, namely SKA continuum and H I observations. For example,

- do UDGs host low-luminosity active galactic nuclei or residual star formation? Deep, high-resolution SKA continuum imaging can reveal weak synchrotron cores or diffuse disk emission that are currently below detection limits, testing whether UDGs are truly quiescent or harbor faint, persistent activity. Thus, distinguish between thermal star-formation-driven emission and weak non-thermal synchrotron signatures from embedded active galaxies, even in the faintest UDGs.
- What is the radio–infrared–optical connection in diffuse systems? SKA’s sensitivity across SKA-Low and SKA-Mid will allow combined spectral studies with Euclid and JWST, constraining dust content, magnetic field strength, and the efficiency of cosmic-ray heating in low-surface-brightness extended diffuse galaxies.
- How does the cluster environment regulate gas stripping and quenching in UDGs? Observations and models agree that richer halos host proportionally more UDGs, but the normalization and scaling exponent ($N_{\text{UDG}} \propto M_{200}^x$) are still being refined (e.g., NUDGEs: $N_{\text{UDG}} \propto M_{200}^{0.78 \pm 0.28}$; Makda et al., 2025). Thus, mapping the spatial correlation between

UDGs and cluster radio halos or relics can test whether environmental processes such as ram-pressure stripping or tidal interactions suppress the radio emission of UDGs.

- Is there a faint-end turnover in the active galactic nuclei luminosity function within clusters? Current luminosity functions may underestimate very low surface-brightness UDGs. Clearly, improved selection tends to increase the number counts of faint sources but uncertainties remain, and different simulations predict different break locations and slopes in the luminosity function. By probing well below current flux density limits, the SKA will quantify the contribution of UDGs and other low-mass systems to the extremely faint active galaxy population, placing constraints on the fraction of galaxies that host a central supermassive black hole in the low-mass regime.
- Can joint SKA continuum and HI spectral-line observations disentangle gas-rich, star-forming UDGs from quenched, dark systems? Combining deep continuum maps with future SKA HI data will directly link the presence (or absence) of cold gas reservoirs to ongoing star formation and feedback, enabling a clear separation between genuinely dark, gas-poor UDGs and those that retain fuel for intermittent star formation within cluster environments.

Together, these investigations will establish whether UDGs represent the low-luminosity, gas-depleted end-state of galaxy evolution or a distinct class of dark-matter-dominated systems with intermittent activity, which is an emerging view from state-of-the-art hydrodynamical simulations (Di Cintio et al., 2017; Sales et al., 2020; Tremmel et al., 2020). Future deep observations with SKA-Mid and SKA-Low will be crucial to distinguish between genuinely quiescent UDGs and those hosting extremely weak active galactic nuclei activity.

Acknowledgments

We thank the anonymous reviewer for many suggestions that greatly improved the manuscript. D.V.L. acknowledges the support of the Department of Atomic Energy, Government of India, under project no. 12-R&D-TFR-5.02-0700. We thank the staff of the GMRT who have made these observations possible. The GMRT is run by the National Centre for Radio Astrophysics of the Tata Institute of Fundamental Research.

References

- H. S. B. Algera et al. In *Advancing Astrophysics with the SKA – II (AASKAII)*. 2026. arXiv search: Report number AASKAII/Algera01.
- F. X. An et al. In *Advancing Astrophysics with the SKA – II (AASKAII)*. 2026. arXiv search: Report number AASKAII/FangxiaAn01.
- R. D. Baldi, A. Capetti, and G. Giovannini. *MNRAS*, 482(2):2294–2304, Jan. 2019. doi: 10.1093/mnras/sty2703.
- C. E. Barbosa et al. *ApJS*, 247(2):46, Apr. 2020. doi: 10.3847/1538-4365/ab7660.
- J. A. Benavides et al. *Nature Astronomy*, 5:1255–1260, Sept. 2021. doi: 10.1038/s41550-021-01458-1.

- J. A. Benavides et al. *MNRAS*, 522(1):1033–1048, June 2023. doi: 10.1093/mnras/stad1053.
- P. Bontempi et al. *MNRAS*, 426(1):588–594, Oct. 2012. doi: 10.1111/j.1365-2966.2012.21786.x.
- T. Bourke et al. *Advancing astrophysics with the square kilometre array (aaska14)*, 2015.
- R. Braun. In *Advancing Astrophysics with the Square Kilometre Array (AASKA14)*, page 174, Apr. 2015. doi: 10.22323/1.215.0174.
- R. Braun et al. Anticipated ska1 science performance. Technical Report SKAO-TEL-0000818, SKA Observatory, 2024. Version 2.0, June 2024.
- T. Carleton et al. *MNRAS*, 485(1):382–395, May 2019. doi: 10.1093/mnras/stz383.
- J. J. Condon et al. *AJ*, 115(5):1693–1716, May 1998. doi: 10.1086/300337.
- A. Di Cintio et al. *MNRAS*, 466(1):L1–L6, Mar. 2017. doi: 10.1093/mnras/1slw210.
- Euclid Collaboration et al. *A&A*, 697:A1, May 2025. doi: 10.1051/0004-6361/202450810.
- B. Fanaroff et al. *MNRAS*, 505(4):6003–6016, Aug. 2021. doi: 10.1093/mnras/stab1540.
- C. Fielder et al. *AJ*, 168(5):212, Nov. 2024. doi: 10.3847/1538-3881/ad74f6.
- J. S. Gannon et al. *MNRAS*, 502(3):3144–3157, Apr. 2021. doi: 10.1093/mnras/stab277.
- J. S. Gannon et al. *MNRAS*, 510(1):946–958, Feb. 2022. doi: 10.1093/mnras/stab3297.
- J. S. Gannon et al. *MNRAS*, 531(1):1856–1869, June 2024. doi: 10.1093/mnras/stae1287.
- J. R. Gott, III, M. S. Vogeley, S. Podariu, and B. Ratra. *ApJ*, 549(1):1–17, Mar. 2001. doi: 10.1086/319055.
- M. J. Hardcastle et al. In *Advancing Astrophysics with the SKA – II (AASKAII)*. 2026. arXiv search: Report number AASKAII/Hardcastle01.
- L. C. Ho, A. V. Filippenko, and W. L. W. Sargent. *ApJ*, 487(2):579–590, Oct. 1997. doi: 10.1086/304642.
- S. R. Janssens et al. *ApJ*, 887(1):92, Dec. 2019. doi: 10.3847/1538-4357/ab536c.
- M. G. Jones et al. *ApJ*, 919(2):72, Oct. 2021. doi: 10.3847/1538-4357/ac0975.
- A. Karunakaran et al. *ApJ*, 975(1):91, Nov. 2024. doi: 10.3847/1538-4357/ad77cf.
- J. Koda, M. Yagi, H. Yamanoi, and Y. Komiyama. *The Astrophysical Journal Letters*, 807:L2, 2015. doi: 10.1088/2041-8205/807/1/L2.
- D. V. Lal and L. C. Ho. *AJ*, 139(3):1089–1105, Mar. 2010. doi: 10.1088/0004-6256/139/3/1089.
- D. V. Lal, P. Shastri, and D. C. Gabuzda. *ApJ*, 731(1):68, Apr. 2011. doi: 10.1088/0004-637X/731/1/68.
- D. V. Lal et al. *ApJ*, 934(2):170, Aug. 2022. doi: 10.3847/1538-4357/ac7a9b.
- J. Li et al. *ApJ*, 955(1):1, Sept. 2023. doi: 10.3847/1538-4357/ace829.

- M. Magliocchetti et al. *arXiv e-prints*, art. arXiv:2511.02970, Nov. 2025. doi: 10.48550/arXiv.2511.02970.
- N. A. Makda, S.-L. Blyth, and R. E. Skelton. *MNRAS*, Oct. 2025. doi: 10.1093/mnras/staf1850.
- G. Martin et al. *MNRAS*, 485(1):796–818, May 2019. doi: 10.1093/mnras/stz356.
- D. Martínez-Delgado et al. *AJ*, 151(4):96, Apr. 2016. doi: 10.3847/0004-6256/151/4/96.
- G. Mazzolari et al. In *Advancing Astrophysics with the SKA – II (AASKAII)*. 2026. arXiv search: Report number AASKAII/Mazzolari01.
- J. C. Mihos et al. *ApJ*, 809(2):L21, Aug. 2015. doi: 10.1088/2041-8205/809/2/L21.
- K. Motiwala et al. *ApJ*, 989(1):86, Aug. 2025. doi: 10.3847/1538-4357/ade9a0.
- E. J. Murphy et al. *ApJ*, 737(2):67, Aug. 2011. doi: 10.1088/0004-637X/737/2/67.
- E. Papastergis, E. A. K. Adams, and A. J. Romanowsky. *A&A*, 601:L10, May 2017. doi: 10.1051/0004-6361/201730795.
- G. L. Pilbratt et al. *A&A*, 518:L1, July 2010. doi: 10.1051/0004-6361/201014759.
- J. Román and I. Trujillo. *MNRAS*, 468(1):703–716, June 2017a. doi: 10.1093/mnras/stx438.
- J. Román and I. Trujillo. *MNRAS*, 468(4):4039–4047, July 2017b. doi: 10.1093/mnras/stx694.
- Y. Rong et al. *MNRAS*, 470(4):4231–4240, Oct. 2017. doi: 10.1093/mnras/stx1440.
- L. V. Sales et al. *MNRAS*, 494(2):1848–1858, May 2020. doi: 10.1093/mnras/staa854.
- T. C. Scott et al. *MNRAS*, 503(3):3953–3964, May 2021. doi: 10.1093/mnras/stab390.
- T. W. Shimwell et al. *A&A*, 659:A1, Mar. 2022. doi: 10.1051/0004-6361/202142484.
- M. F. Struble. *MNRAS*, 473(4):4686–4691, Feb. 2018. doi: 10.1093/mnras/stx1785.
- M. Tremmel et al. *MNRAS*, 497(3):2786–2810, Sept. 2020. doi: 10.1093/mnras/staa2015.
- J. S. Ulvestad and L. C. Ho. *ApJ*, 562(2):L133–L136, Dec. 2001. doi: 10.1086/338254.
- J. S. Ulvestad and L. C. Ho. *ApJ*, 581(2):925–931, Dec. 2002. doi: 10.1086/344442.
- R. F. J. van der Burg, A. Muzzin, and H. Hoekstra. *A&A*, 590:A20, May 2016. doi: 10.1051/0004-6361/201628222.
- R. F. J. van der Burg et al. *A&A*, 607:A79, Nov. 2017. doi: 10.1051/0004-6361/201731335.
- P. van Dokkum et al. *Nature*, 605(7910):435–439, May 2022. doi: 10.1038/s41586-022-04665-6.
- P. G. van Dokkum et al. *ApJ*, 798(2):L45, Jan. 2015. doi: 10.1088/2041-8205/798/2/L45.
- J. D. Van Nest et al. *ApJ*, 926(1):92, Feb. 2022. doi: 10.3847/1538-4357/ac43b7.
- R. L. White et al. *ApJ*, 654(1):99–114, Jan. 2007. doi: 10.1086/507700.

M. Yagi, J. Koda, Y. Komiyama, and H. Yamanai. *ApJS*, 225(1):11, July 2016. doi: 10.3847/0067-0049/225/1/11.

D. Zaritsky et al. *ApJS*, 240(1):1, Jan. 2019. doi: 10.3847/1538-4365/aafe9.

D. Zaritsky et al. *ApJS*, 257(2):60, Dec. 2021. doi: 10.3847/1538-4365/ac2607.

H. Zheng, S. Liao, L. Gao, and F. Jiang. *arXiv e-prints*, art. arXiv:2504.14973, Apr. 2025. doi: 10.48550/arXiv.2504.14973.

Y. Zheng et al. *ApJ*, 960(1):55, Jan. 2024. doi: 10.3847/1538-4357/acfe6b.